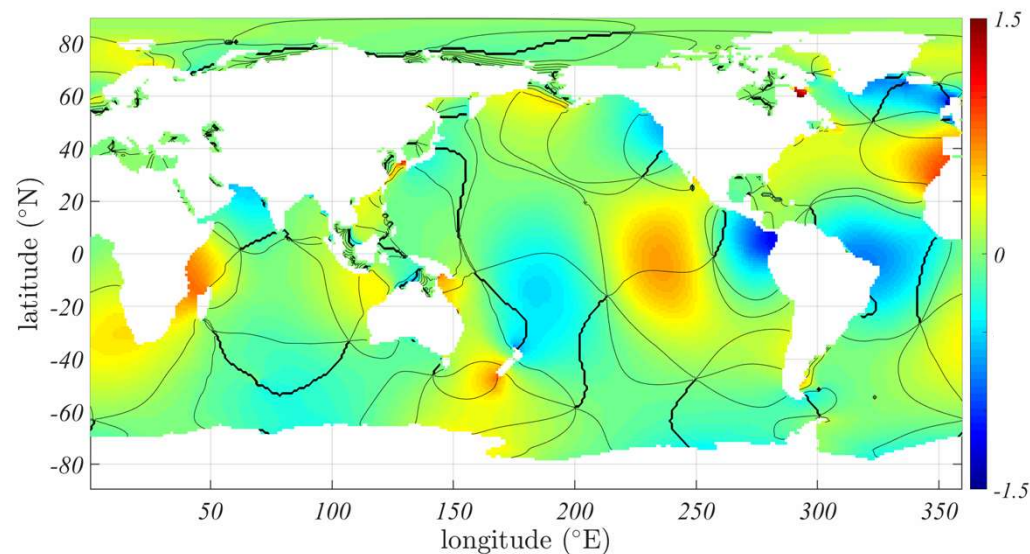


# Lecture 6: Shallow water dynamics and gravity waves

Callum J. Shakespeare

*Fellow, Climate and Fluid Physics, ANU*



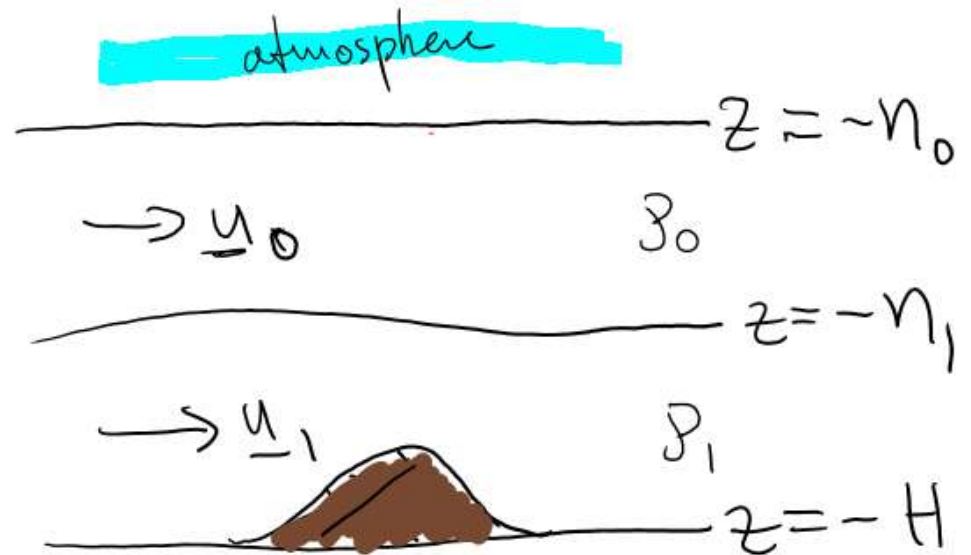
# Derivation of 2-layer shallow water equations

Hydrostatic momentum equations

$$\frac{Du}{Dt} - fv = -\frac{1}{\rho} \frac{\partial p}{\partial x}$$

$$\frac{Dv}{Dt} + fu = -\frac{1}{\rho} \frac{\partial p}{\partial y}$$

$$\frac{\partial p}{\partial z} = \rho g$$



$$p_0 = p_0 g (-n_0 - z)$$

$$p_1 = p_0 g (-n_0 + n_1) + p_1 g (-n_1 - z)$$

# Derivation of 2-layer shallow water equations

$$\frac{\partial p}{\partial z} = \rho g$$



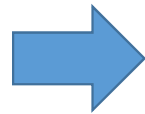
$$P_0 = \rho_0 g (-\eta_0 - z)$$

$$P_1 = \rho_0 g (-\eta_0 + \eta_1) + \rho_1 g (-\eta_1 - z)$$



$$\frac{Du}{Dt} - fv = -\frac{1}{\rho} \frac{\partial p}{\partial x}$$

$$\frac{Dv}{Dt} + fu = -\frac{1}{\rho} \frac{\partial p}{\partial y}$$

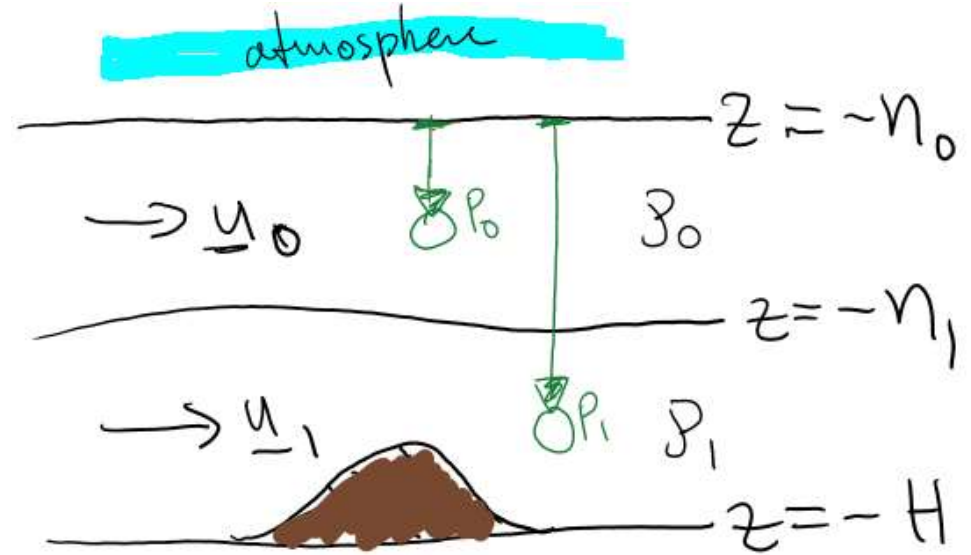


Layer 0

$$\frac{Du_0}{Dt} - fv_0 = g \frac{\partial \eta_0}{\partial x}$$

Layer 1

$$\begin{aligned} \frac{Du_1}{Dt} - fv_1 &= g \frac{\partial(\eta_0 - \eta_1)}{\partial x} + g \frac{\rho_1}{\rho_0} \frac{\partial \eta_1}{\partial x} \\ &= g \frac{\partial \eta_0}{\partial x} + \boxed{g' \frac{\partial \eta_1}{\partial x}} \end{aligned}$$



# Derivation of 2-layer shallow water equations

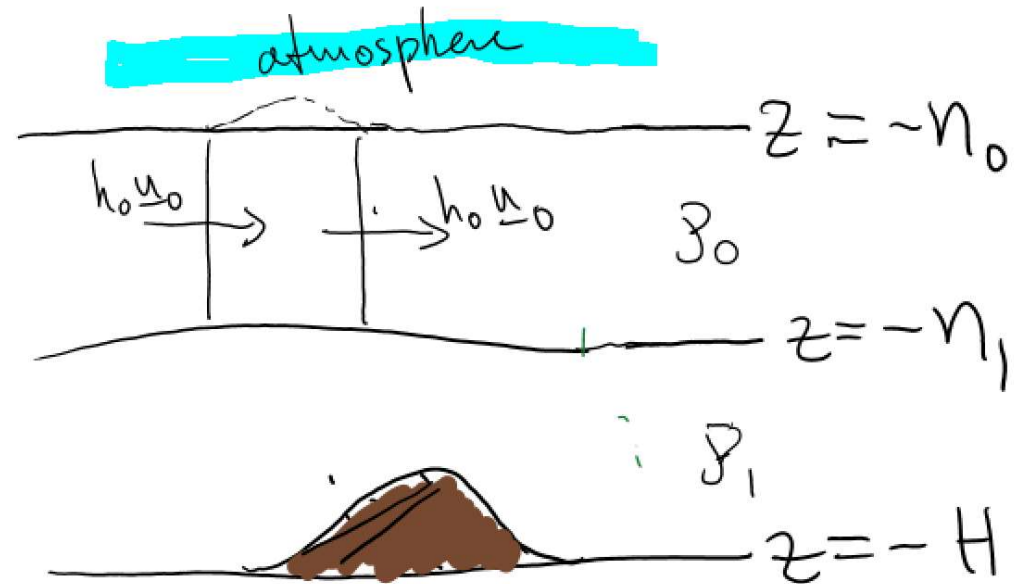
Conservation of mass

$$\begin{aligned}
 h_0 &= \eta_1 - \eta_0 \\
 \frac{\partial h_0}{\partial t} &= -\nabla \cdot (\text{volume flux}) \\
 &= -\nabla \cdot (h_0 \underline{u}_0) \\
 &= -h_0 \nabla \cdot \underline{u}_0 - \underline{u}_0 \cdot \nabla h
 \end{aligned}$$



$$\frac{Dh_0}{Dt} = -h_0 \nabla \cdot \underline{u}_0$$

Layer 1 similar but  $h_1 = H - \eta_1$



## 2-layer case

$$\frac{Du_0}{Dt} - fv_0 = g \frac{\partial \eta_0}{\partial x}$$

$$\frac{Dv_0}{Dt} + fu_0 = g \frac{\partial \eta_0}{\partial y}$$

$$\frac{Dh_0}{Dt} = -h_0 \nabla \cdot \mathbf{u}_0$$

$$h_0 = \eta_1 - \eta_0$$

$$\frac{Du_1}{Dt} - fv_1 = g \frac{\partial \eta_0}{\partial x} + g' \frac{\partial \eta_1}{\partial x}$$

$$\frac{Dv_1}{Dt} + fu_1 = g \frac{\partial \eta_0}{\partial y} + g' \frac{\partial \eta_1}{\partial y}$$

$$\frac{Dh_1}{Dt} = -h_1 \nabla \cdot \mathbf{u}_1$$

$$h_1 = H - \eta_1$$



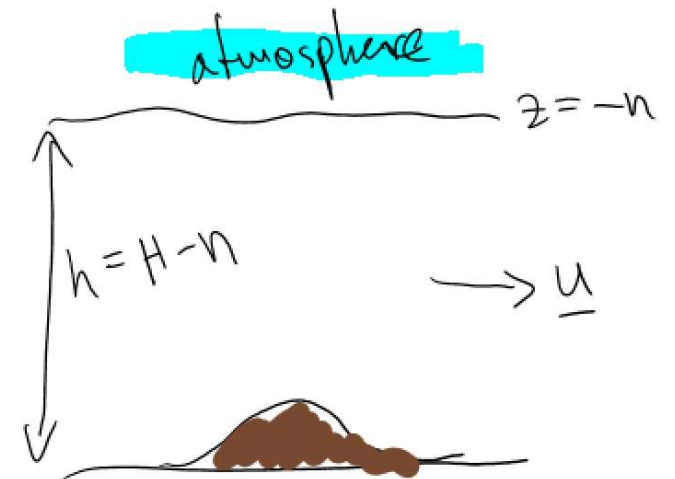
## 1-layer case

$$\frac{Du}{Dt} - fv = g \frac{\partial \eta}{\partial x}$$

$$\frac{Dv}{Dt} + fu = g \frac{\partial \eta}{\partial y}$$

$$\frac{Dh}{Dt} = -h \nabla \cdot \mathbf{u}$$

$$h = H - \eta$$



# Linearisation and wave solutions

Non-linear equations

$$\frac{Du}{Dt} - fv = g \frac{\partial \eta}{\partial x}$$

$$\frac{Dv}{Dt} + fu = g \frac{\partial \eta}{\partial y}$$

$$\frac{Dh}{Dt} = -h \nabla \cdot \mathbf{u}$$

$$h = H - \eta$$

Linearised equations

$$\frac{\partial u}{\partial t} - fv = g \frac{\partial \eta}{\partial x}$$

$$\frac{\partial v}{\partial t} + fu = g \frac{\partial \eta}{\partial y}$$

$$-\frac{\partial \eta}{\partial t} + \nabla \cdot (\underline{u} H) = 0$$

Assumptions

$$\left( \underline{u} \cdot \nabla \ll \frac{\partial}{\partial t} \Rightarrow \frac{u}{\omega L} \ll 1 \right)$$

$$(\eta \ll H)$$

# Linearisation

Non-linear equations

$$\frac{Du}{Dt} - fv = g \frac{\partial \eta}{\partial x}$$

$$\frac{Dv}{Dt} + fu = g \frac{\partial \eta}{\partial y}$$

$$\frac{Dh}{Dt} = -h \nabla \cdot \mathbf{u}$$

$$h = H - \eta$$

Linearised equations

$$\frac{\partial u}{\partial t} - f v = g \frac{\partial \eta}{\partial x}$$

$$\frac{\partial v}{\partial t} + f u = g \frac{\partial \eta}{\partial y}$$

$$-\frac{\partial \eta}{\partial t} + \nabla \cdot (\mathbf{u} H) = 0$$

Assumptions

$$\left( \underline{u} \cdot \nabla \ll \frac{\partial}{\partial t} \Rightarrow \frac{u}{\omega L} \ll 1 \right)$$

$$(\eta \ll H)$$

$$\text{let } \underline{U} = \underline{u} H$$

$$\frac{\partial U}{\partial t} - f V = g H \frac{\partial \eta}{\partial x}$$

$$\frac{\partial V}{\partial t} + f U = g H \frac{\partial \eta}{\partial y}$$

$$-\frac{\partial \eta}{\partial t} + \nabla \cdot \underline{U} = 0$$

Depth  
integrated  
linearised  
equations

# Wave solutions

$$\begin{aligned}\frac{\partial u}{\partial t} - fV &= gH \frac{\partial \eta}{\partial x} \\ \frac{\partial v}{\partial t} + fU &= gH \frac{\partial \eta}{\partial y} \\ -\frac{\partial \eta}{\partial t} + \nabla \cdot \underline{U} &= 0\end{aligned}$$

Let  $\phi = \hat{\phi} e^{i(kx + ly - \omega t)}$

$$\begin{aligned}\Rightarrow -i\omega \hat{u} - f\hat{v} &= ikgH \hat{\eta} \\ -i\omega \hat{v} + f\hat{u} &= ilgH \hat{\eta} \\ +i\omega \hat{\eta} + ik\hat{u} + il\hat{v} &= 0\end{aligned}$$



## Dispersion Relation

$$\omega^2 = f^2 + (k^2 + l^2)gH$$

$$\omega = \pm \sqrt{f^2 + K^2 gH}$$

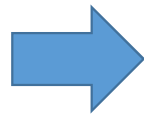
Barotropic gravity waves (BGWs)



# How do we generate BGWs?

- Full depth (barotropic) forcing = body force
- E.g. horizontal variations in gravity due to the moon = tidal forces!
- Usually represented in terms of an equilibrium tide = height perturbation that would exist if no flow:  $\bar{\eta} = A \cos(\omega_t t + 2\phi)$
- Flow is associated with departures from  $\bar{\eta}$
- Let  $\eta = \bar{\eta} + \eta'$

$$\begin{aligned} \frac{\partial u}{\partial t} - fV &= gH \frac{\partial \eta}{\partial x} \\ \frac{\partial v}{\partial t} + fU &= gH \frac{\partial \eta}{\partial y} \\ - \frac{\partial \eta}{\partial t} + \nabla \cdot \underline{U} &= 0 \end{aligned}$$



$$\begin{aligned} \frac{\partial u}{\partial t} - fV &= gH \frac{\partial \eta'}{\partial x} \\ \frac{\partial v}{\partial t} + fU &= gH \frac{\partial \eta'}{\partial y} \\ - \frac{\partial \eta'}{\partial t} + \nabla \cdot \underline{U} &= \boxed{+ \frac{\partial \bar{\eta}}{\partial t}} \end{aligned}$$

Tidal forcing

These are the Tidal Equations  
or if on a sphere "Laplace Tidal Equations"

# Properties of the tides

- Forced at a particular frequency = 12.4 hour period

$$\omega = \pm \sqrt{f^2 + K^2 g H}$$

- $f \sim \frac{10^{-4}}{s}$ ,  $\omega \sim 1.4 \times \frac{10^{-4}}{s}$ ,  $g \sim 10 \frac{m}{s^2}$ ,  $H \sim 4000 m$
- What is the value of K?
- $\frac{1}{K} \sim 2000 km$

# Properties of the tides

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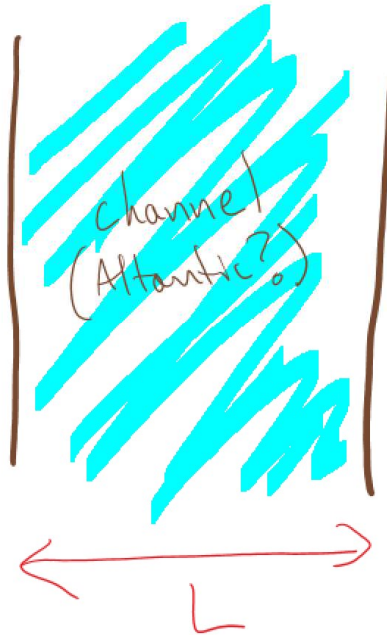
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- $\frac{1}{K} \sim 2000 km$

How is the size of my ocean basin related to the amplitude of the tide????



# Properties of the tides

How is the size of my ocean basin related to the amplitude of the tide????



- Consider a simplified channel
- Infinitely long in  $y$
- Width  $L$  in  $x$

# Properties of the tides

How is the size of my ocean basin related to the amplitude of the tide????



$$\frac{\partial u}{\partial t} - fV = -gH \frac{\partial \eta'}{\partial x}$$

$$\frac{\partial v}{\partial t} + fU = 0$$

$$\frac{\partial \eta'}{\partial t} + \frac{\partial \eta'}{\partial x} = -\frac{\partial \bar{\eta}}{\partial t}$$

$$\left. \begin{array}{l} \frac{\partial \eta'}{\partial t} - fV = -gH \frac{\partial \eta'}{\partial x} \\ \frac{\partial v}{\partial t} + fU = 0 \end{array} \right\} \bar{\eta} = A \cos\left(\omega t + \frac{2x}{C}\right)$$

$$C = 2\pi R_E \cos\theta$$

Circumference of Earth

$$\bar{\eta} \approx A \cos(\omega t) - \frac{4Ax}{C} \sin(\omega t) \quad \left(\frac{2x}{C} \ll \omega t\right)$$

# Properties of the tides

How is the size of my ocean basin related to the amplitude of the tide????



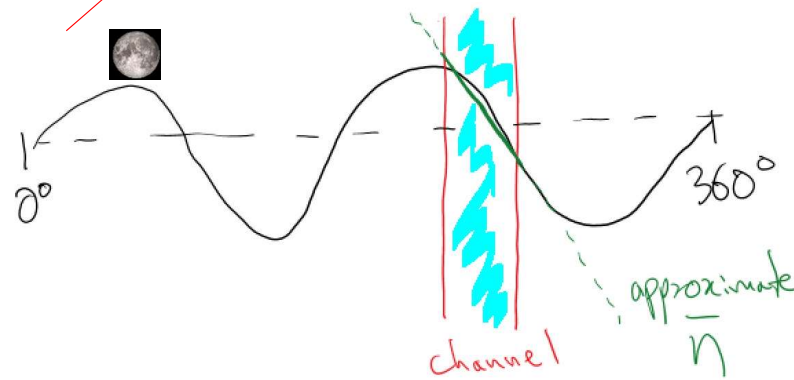
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# Properties of the tides

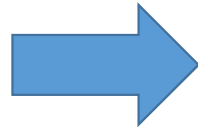
How is the size of my ocean basin related to the amplitude of the tide???

$$\text{Let: } \eta' = \text{Im}(\hat{\eta}' e^{i\omega t})$$

$$i\omega \hat{u} - f\hat{v} = -gH \frac{\partial \hat{\eta}'}{\partial x}$$

$$i\omega \hat{v} + f\hat{u} = 0 \Rightarrow \hat{v} = -\frac{f}{\omega} \hat{u}$$

$$i\omega(\hat{\eta}' - \hat{\eta}) + \frac{\partial \hat{u}}{\partial x} = 0$$



$$\text{Thus: } i\omega \hat{u} + \frac{f^2}{i\omega} \hat{u} = -gH \frac{\partial \hat{\eta}'}{\partial x}$$

$$\Rightarrow \hat{u} = -\frac{i\omega}{f^2 - \omega^2} gH \frac{\partial \hat{\eta}'}{\partial x}$$

$$\Rightarrow i\omega(\hat{\eta}' - A) - \frac{i\omega gH}{f^2 - \omega^2} \frac{\partial^2 \hat{\eta}'}{\partial x^2} = 0$$

$$\boxed{\hat{\eta}' - \frac{gH}{f^2 - \omega^2} \frac{\partial^2 \hat{\eta}'}{\partial x^2} = -\frac{kAx}{c}}$$

$$\text{where } \hat{u} \sim \frac{\partial \hat{\eta}'}{\partial x} = 0 \text{ at } x=0, L$$

# Properties of the tides

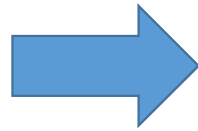
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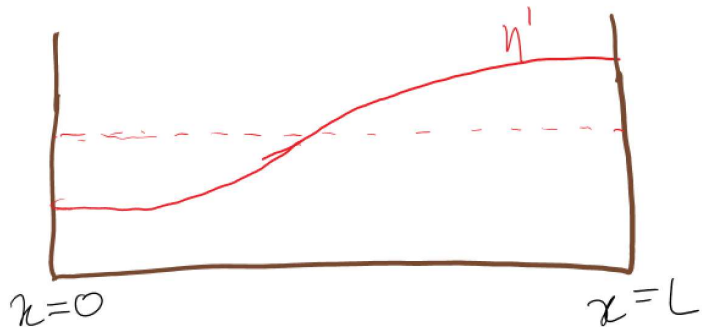
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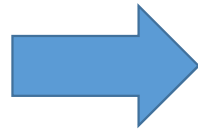
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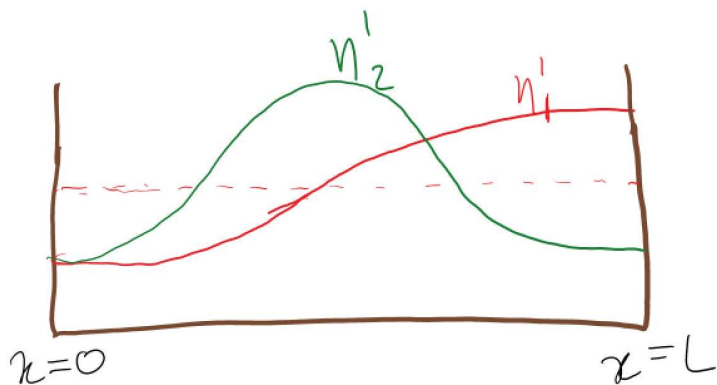
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$$\text{Let } \hat{\eta}' = \sum B_n \cos\left(\frac{n\pi x}{L}\right)$$



# Properties of the tides

How is the size of my ocean basin related to the amplitude of the tide????

$$\text{Let } \hat{h}' = \sum B_n \cos\left(\frac{n\pi x}{L}\right)$$

$$\hat{h}' - \frac{gH}{f^2 - \omega^2} \frac{\partial^2 \hat{h}'}{\partial x^2} = -\frac{kAx}{C}$$

$$\sum_n B_n \cos\left(\frac{n\pi x}{L}\right) - \frac{gH}{f^2 - \omega^2} \frac{\partial^2}{\partial x^2} B_n \cos\left(\frac{n\pi x}{L}\right) = -\frac{4A}{C} \sum_n C_n \cos\left(\frac{n\pi x}{L}\right)$$

where  $x = \sum_n C_n \cos\left(\frac{n\pi x}{L}\right)$

# Properties of the tides

How is the size of my ocean basin related to the amplitude of the tide????

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where  $x = \sum_n C_n \cos\left(\frac{n\pi x}{L}\right) \times \cos\left(\frac{m\pi x}{L}\right) \int_0^L dx$

# Properties of the tides

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where  $x = \sum_n C_n \cos\left(\frac{n\pi x}{L}\right) \times \cos\left(\frac{m\pi x}{L}\right) \int_0^L dx$

$$C_n = \frac{\int_0^L \cos\left(\frac{n\pi x}{L}\right) x dx}{\int_0^L \cos^2\left(\frac{n\pi x}{L}\right) dx}$$

$$= \begin{cases} 0 & n \text{ even} \\ \frac{2L}{n^2\pi^2} & n \text{ odd} \end{cases}$$

$$\hat{h}' = -\frac{4A}{C} \sum_{n=1,3,\dots}^{\infty} \frac{2L}{n^2\pi^2} \frac{1}{1 - \frac{gHn^2\pi^2}{L^2(\omega^2 - f^2)}} \cos\left(\frac{n\pi x}{L}\right)$$

Mode  $n=3$  is ~10% of mode  $n=1$ ... so ignore it.

# Properties of the tides

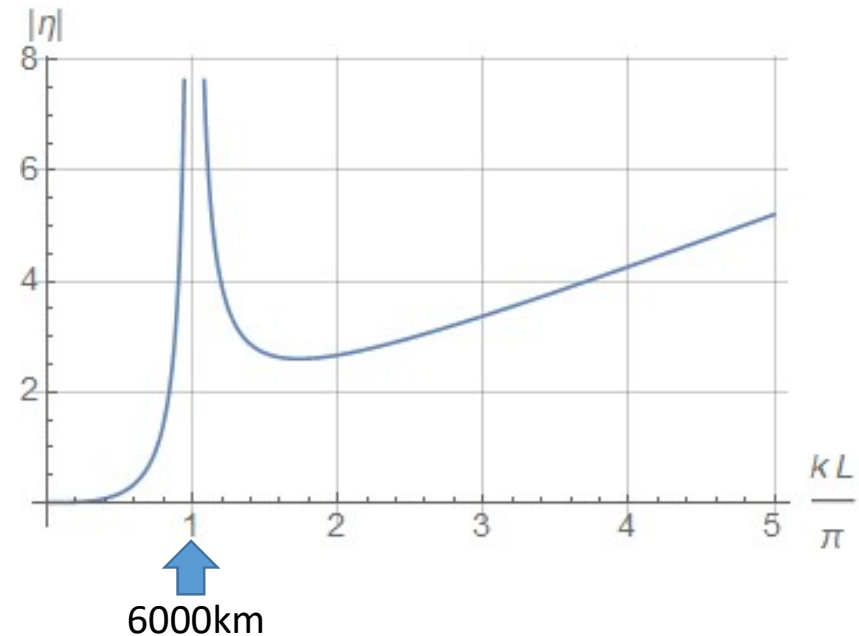
How is the size of my ocean basin related to the amplitude of the tide???

$$\hat{\eta} \approx -\frac{4A}{c} \frac{2L}{\pi^2} \frac{1}{1 - \left(\frac{\pi c}{KL}\right)^2} \cos\left(\frac{\pi x}{L}\right)$$

Here we have used the previous definition

$$\omega = \pm \sqrt{f^2 + K^2 g H}$$

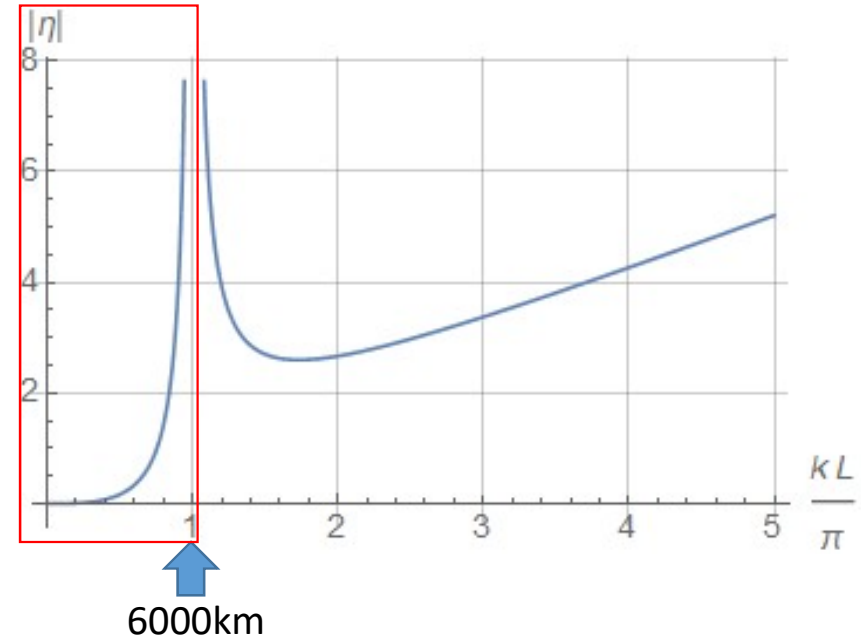
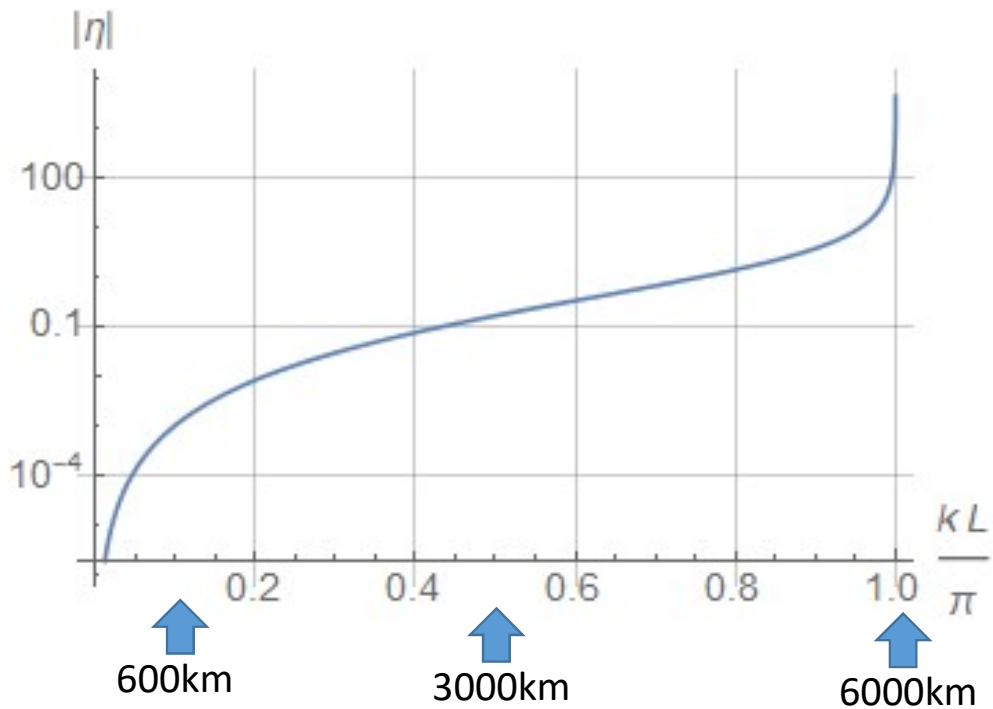
where  $\frac{1}{K} \sim 2000 \text{ km}$



# Properties of the tides

How is the size of my ocean basin related to the amplitude of the tide???

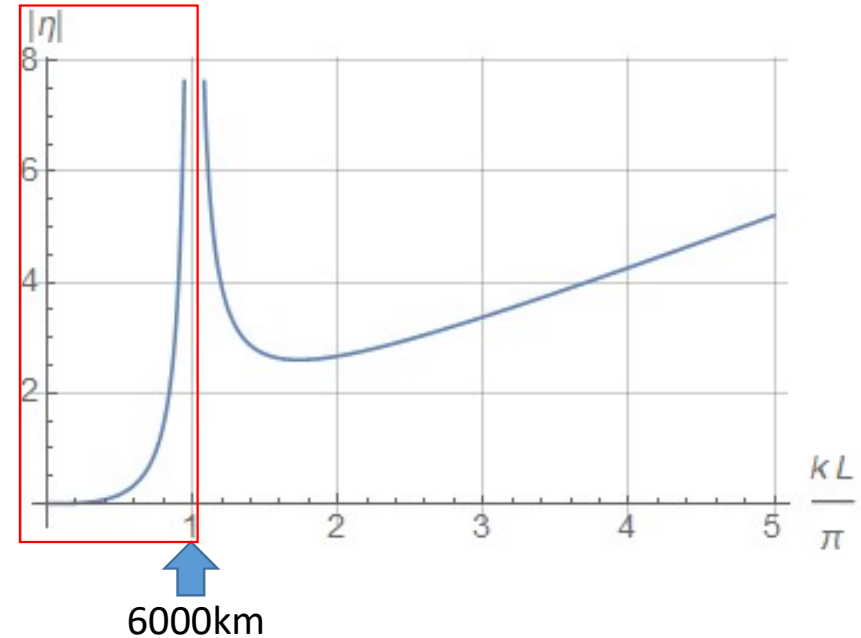
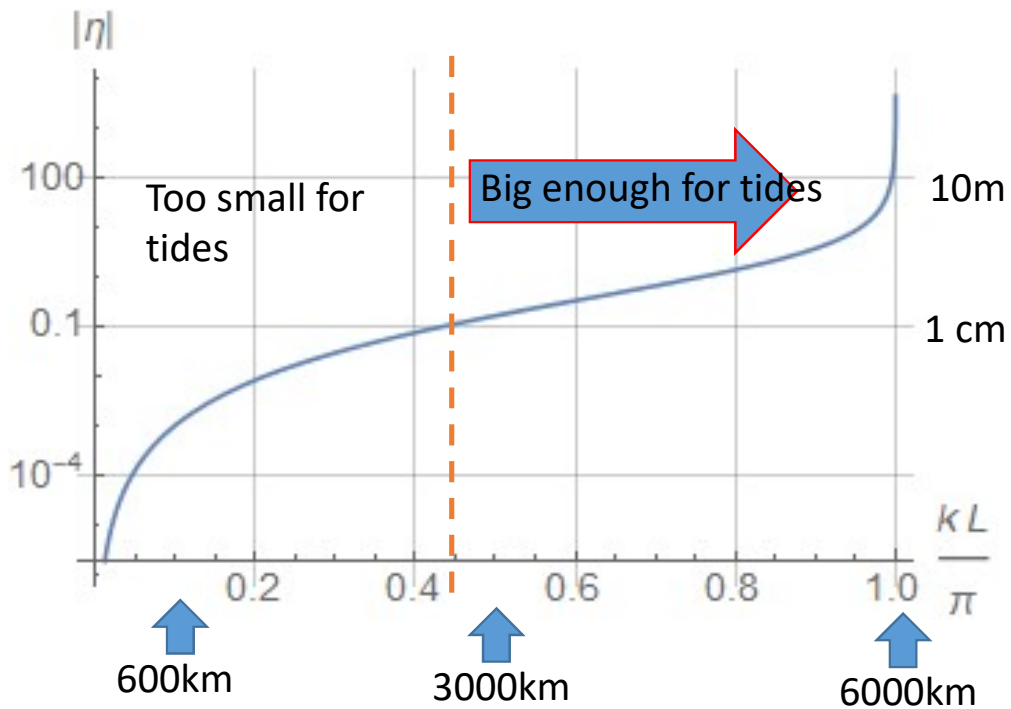
$$\hat{\eta} \approx -\frac{4A}{c} \frac{2L}{\pi^2} \frac{1}{1 - \left(\frac{\pi c}{kL}\right)^2} \cos\left(\frac{\pi x}{L}\right)$$



# Properties of the tides

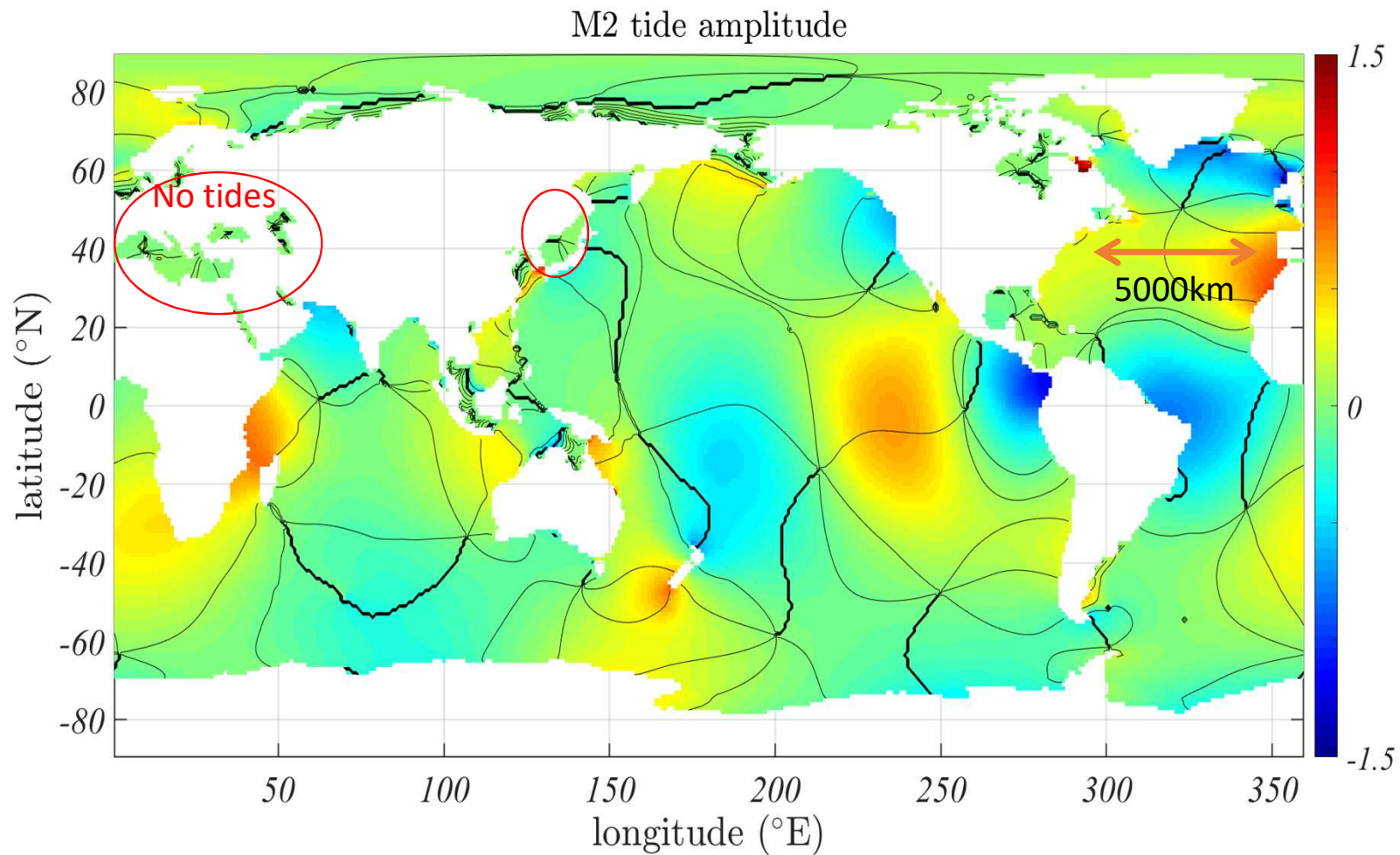
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# Comparison

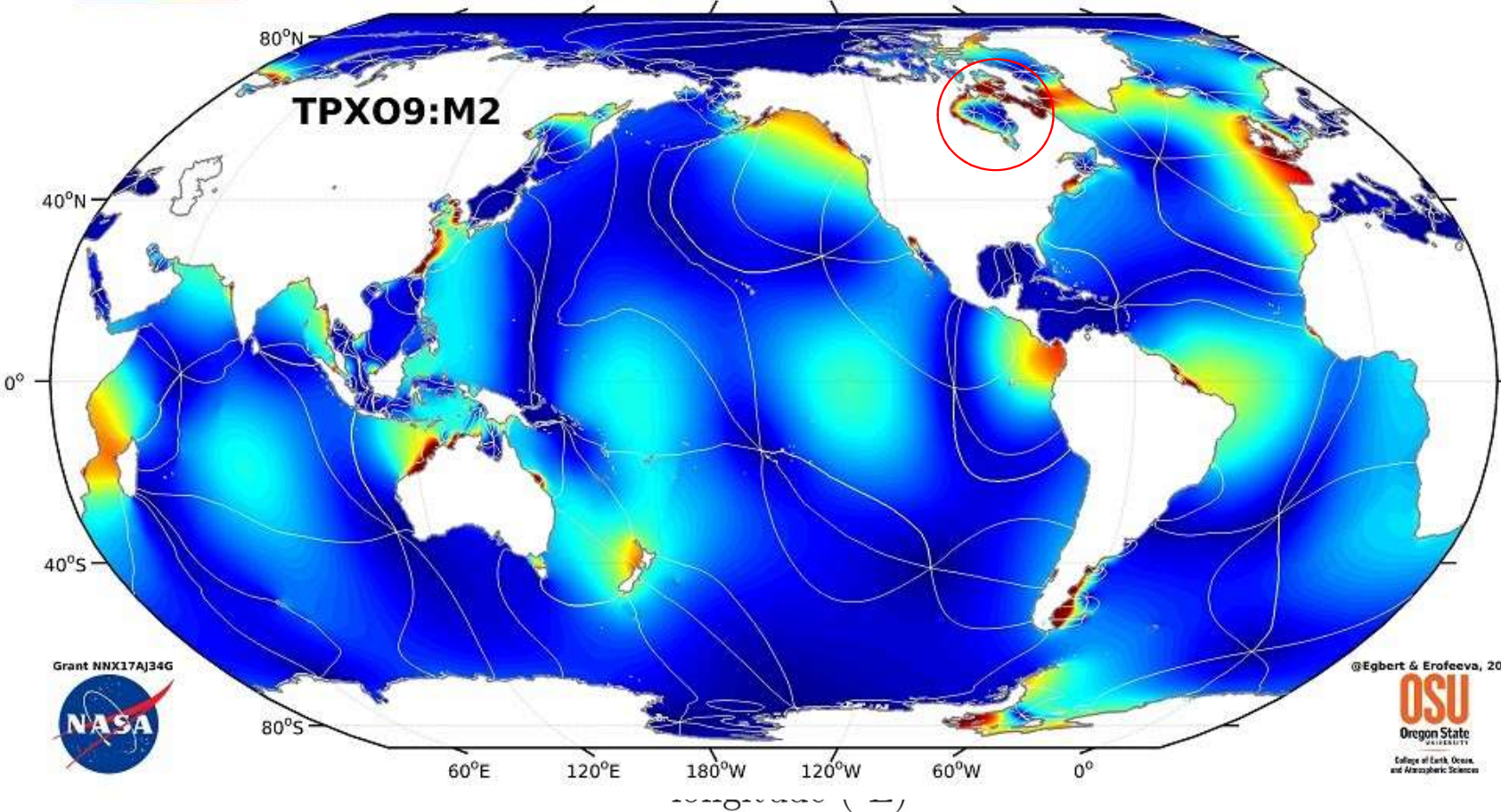
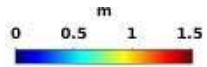
How is the size of my ocean basin related to the amplitude of the tide????





# Comparison

How is the size of my ocean basin related to the amplitude of the tide????



Hudson Bay  
270 m deep max

$L > 500\text{km}$  for tides

<https://www.tpxo.net/global>

# Comparison with Atlantic basin over time

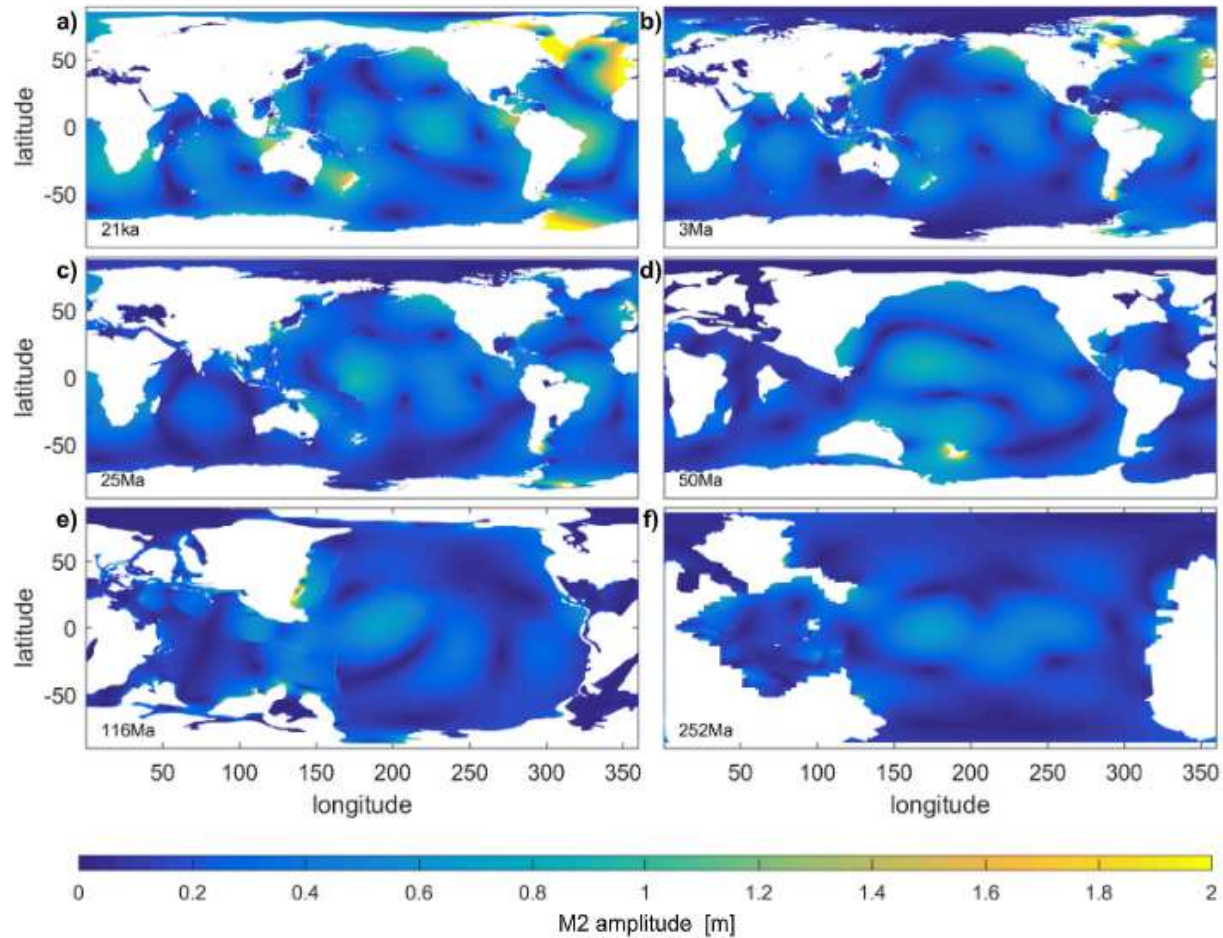


Fig. 2. Shown are the M<sub>2</sub> tidal amplitudes for the LGM (a), Pliocene (b), Miocene (c), Eocene (d), Cretaceous (e) and Permian-Triassic (f).

From Green et al., 2017

# Comparison with Atlantic basin over time

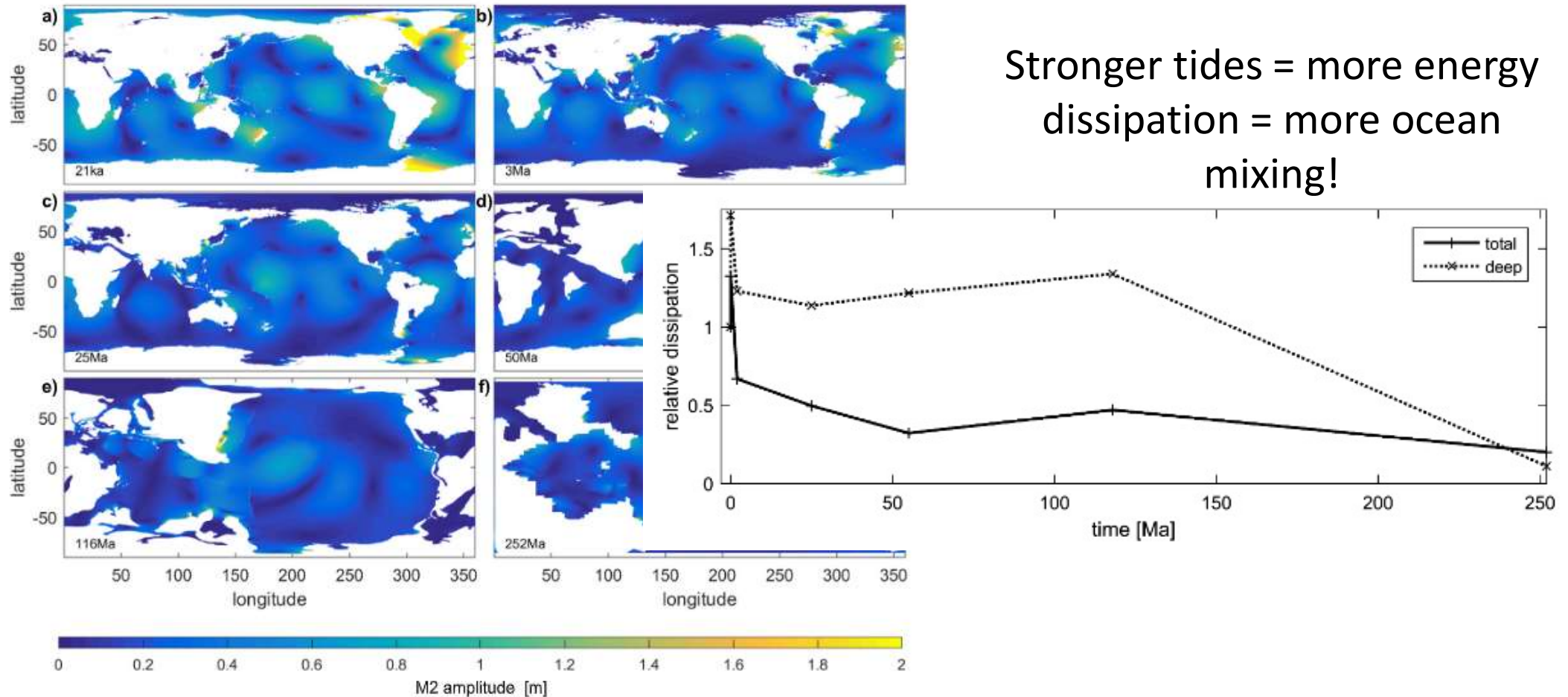
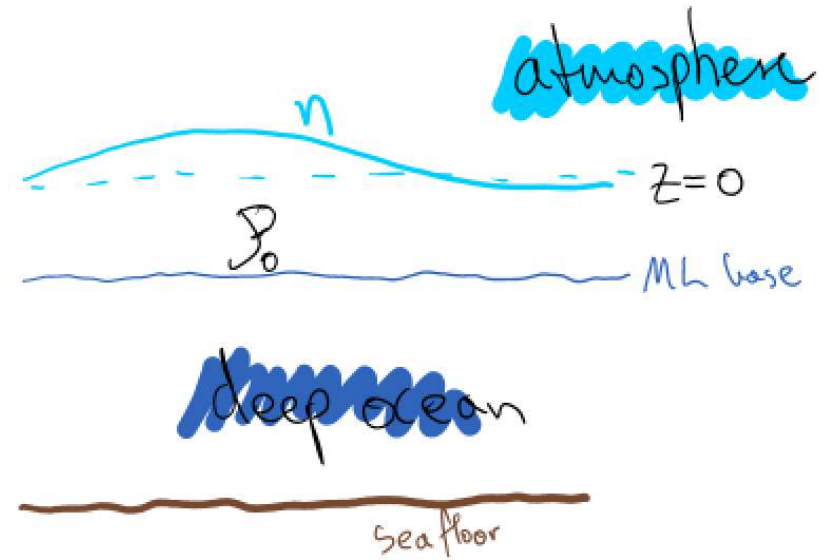


Fig. 2. Shown are the  $M_2$  tidal amplitudes for the LGM (a), Pliocene (b), Miocene (c), Eocene (d), Cretaceous (e) and Permian-Triassic (f).

From Green et al., 2017

# BGWs versus surface gravity waves (SGWs)



# BGWs versus surface gravity waves (SGWs)

Kinematic boundary  
condition

Surface equations

$$w = \frac{Dn}{Dt}$$

Uniform density, **non-hydrostatic** Navier Stokes

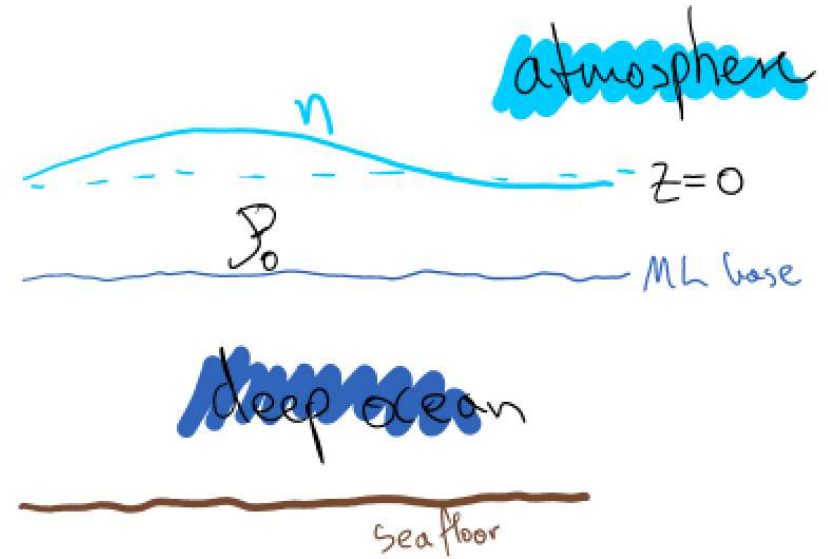
Interior equations

$$\frac{Du}{Dt} = -\frac{1}{\rho_0} \frac{\partial p}{\partial x}$$

$$\frac{Dv}{Dt} = -\frac{1}{\rho_0} \frac{\partial p}{\partial y}$$

$$\frac{Dw}{Dt} = -\frac{1}{\rho_0} \frac{\partial p}{\partial z} + g$$

$$\nabla \cdot \underline{u} = 0.$$



# BGWs versus surface gravity waves (SGWs)

Linearised surface  
( $z=0$ )

$$w = \frac{\partial \eta}{\partial t}$$

Linearised interior

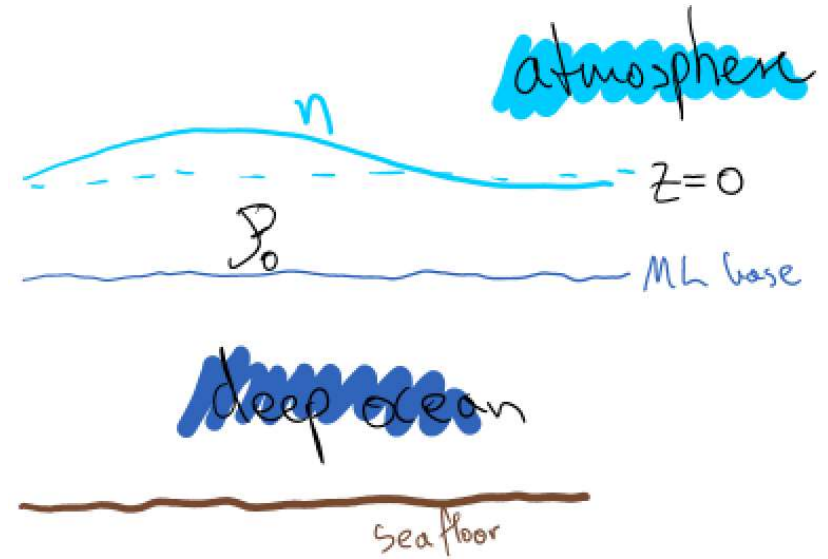
$$\text{let: } p = \rho_0 g z + p'$$

$$\frac{\partial u}{\partial t} = -\frac{1}{\rho_0} \frac{\partial p'}{\partial x}$$

$$\frac{\partial v}{\partial t} = -\frac{1}{\rho_0} \frac{\partial p'}{\partial y}$$

$$\frac{\partial w}{\partial t} = -\frac{1}{\rho_0} \frac{\partial p'}{\partial z}$$

$$\frac{\partial u}{\partial x} + \frac{\partial v}{\partial y} + \frac{\partial w}{\partial z} = 0$$



# BGWs versus surface gravity waves (SGWs)

Linearised surface  
(z=0)

$$w = \frac{\partial \eta}{\partial t}$$



$$\frac{\partial^2 \eta}{\partial t^2} = \frac{\partial w}{\partial z} = -\frac{1}{\rho_0} \frac{\partial p'}{\partial z} \Big|_{z=0}$$

Linearised interior

Let:  $p = \rho_0 g z + p'$

$$\frac{\partial u}{\partial t} = -\frac{1}{\rho_0} \frac{\partial p'}{\partial x}$$

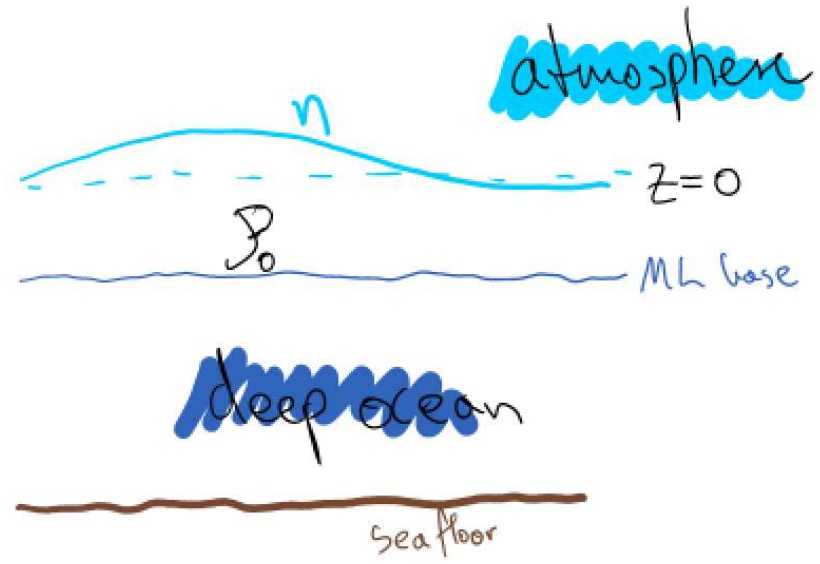
$$\frac{\partial v}{\partial t} = -\frac{1}{\rho_0} \frac{\partial p'}{\partial y}$$

$$\frac{\partial w}{\partial t} = -\frac{1}{\rho_0} \frac{\partial p'}{\partial z}$$

$$\frac{\partial u}{\partial x} + \frac{\partial v}{\partial y} + \frac{\partial w}{\partial z} = 0$$



$$\frac{\partial}{\partial t} (\nabla^2 p') = 0$$



# BGWs versus surface gravity waves (SGWs)

Linearised surface  
(z=0)

$$w = \frac{\partial \eta}{\partial t}$$



$$\frac{\partial^2 \eta}{\partial t^2} = \frac{\partial w}{\partial z} = -\frac{1}{\rho_0} \frac{\partial P}{\partial z} \Big|_{z=0}$$

Linearised interior

Let:  $P = \rho_0 g z + P'$

$$\frac{\partial u}{\partial t} = -\frac{1}{\rho_0} \frac{\partial P'}{\partial x}$$

$$\frac{\partial v}{\partial t} = -\frac{1}{\rho_0} \frac{\partial P'}{\partial y}$$

$$\frac{\partial w}{\partial t} = -\frac{1}{\rho_0} \frac{\partial P'}{\partial z}$$

$$\frac{\partial u}{\partial x} + \frac{\partial v}{\partial y} + \frac{\partial w}{\partial z} = 0$$

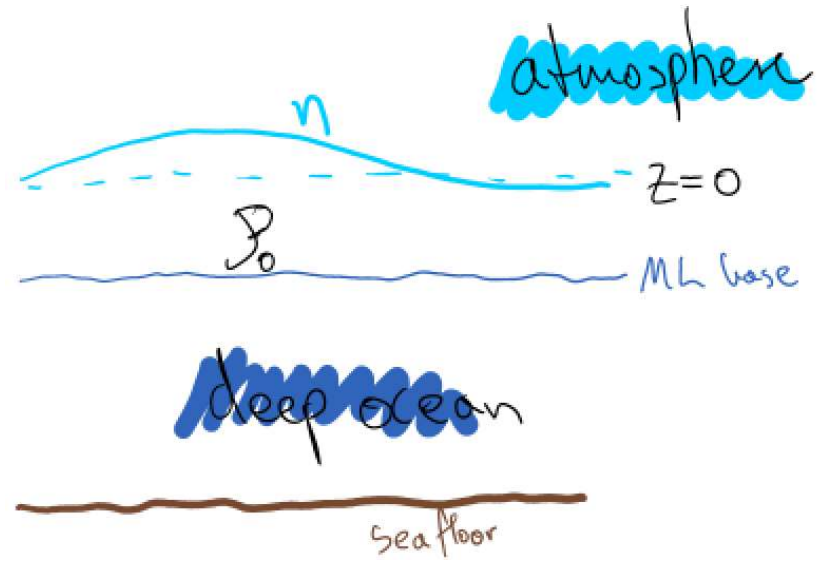


$$\frac{\partial}{\partial t} (\nabla^2 P') = 0$$

Let  $P' \sim \hat{P}' e^{i(kx+ly-wt)}$

Then  $(-(k^2+l^2) + \frac{\partial^2}{\partial z^2}) \hat{P}' = 0$

$\Rightarrow \hat{P}' = \hat{P}'_0 e^{-Kz}$   
where  $K = \sqrt{k^2+l^2}$



Wave amplitude decays with depth, according to horizontal wavelength



# BGWs versus surface gravity waves (SGWs)

Linearised surface  
(z=0)

$$w = \frac{\partial \eta}{\partial t}$$



$$\frac{\partial^2 \eta}{\partial t^2} = \frac{\partial w}{\partial t} = -\frac{1}{\rho_0} \frac{\partial P}{\partial z} \Big|_{z=0}$$

Linearised interior

Let:  $P = \rho_0 g z + P'$

$$\frac{\partial u}{\partial t} = -\frac{1}{\rho_0} \frac{\partial P'}{\partial z}$$

$$\frac{\partial v}{\partial t} = -\frac{1}{\rho_0} \frac{\partial P'}{\partial y}$$

$$\frac{\partial w}{\partial t} = -\frac{1}{\rho_0} \frac{\partial P'}{\partial z}$$

$$\frac{\partial u}{\partial x} + \frac{\partial v}{\partial y} + \frac{\partial w}{\partial z} = 0$$



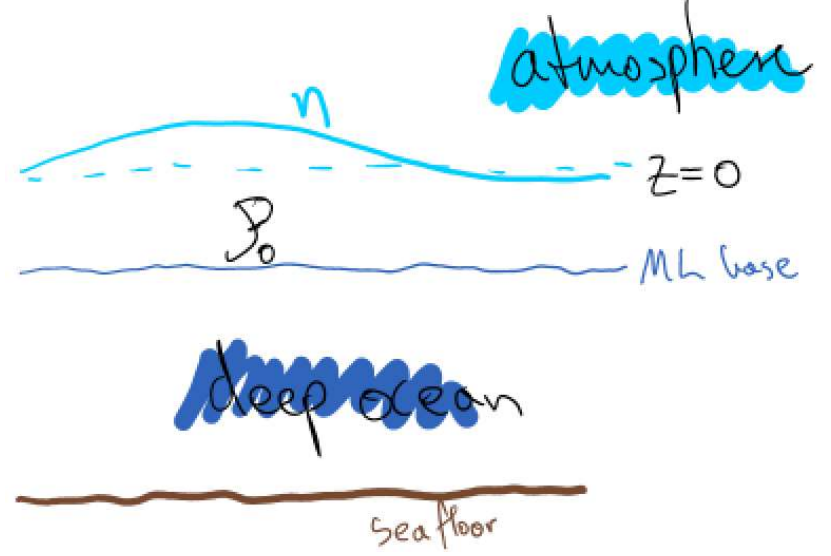
$$\frac{\partial}{\partial t} (\nabla^2 P') = 0$$

Let  $P' \sim \hat{P}' e^{i(kx+ly-wt)}$

Then  $(-(k^2+l^2) + \frac{\partial^2}{\partial z^2}) \hat{P}' = 0$

$$\Rightarrow \hat{P}' = \hat{P}'_0 e^{-Kz}$$

where  $K = \sqrt{k^2+l^2}$



At the surface the total pressure must be constant to match atmospheric pressure, thus

$$P = P' + \rho_0 g \eta = 0 \text{ at surface}$$

$$\Rightarrow \hat{P}'_0 = -\rho_0 g \eta$$

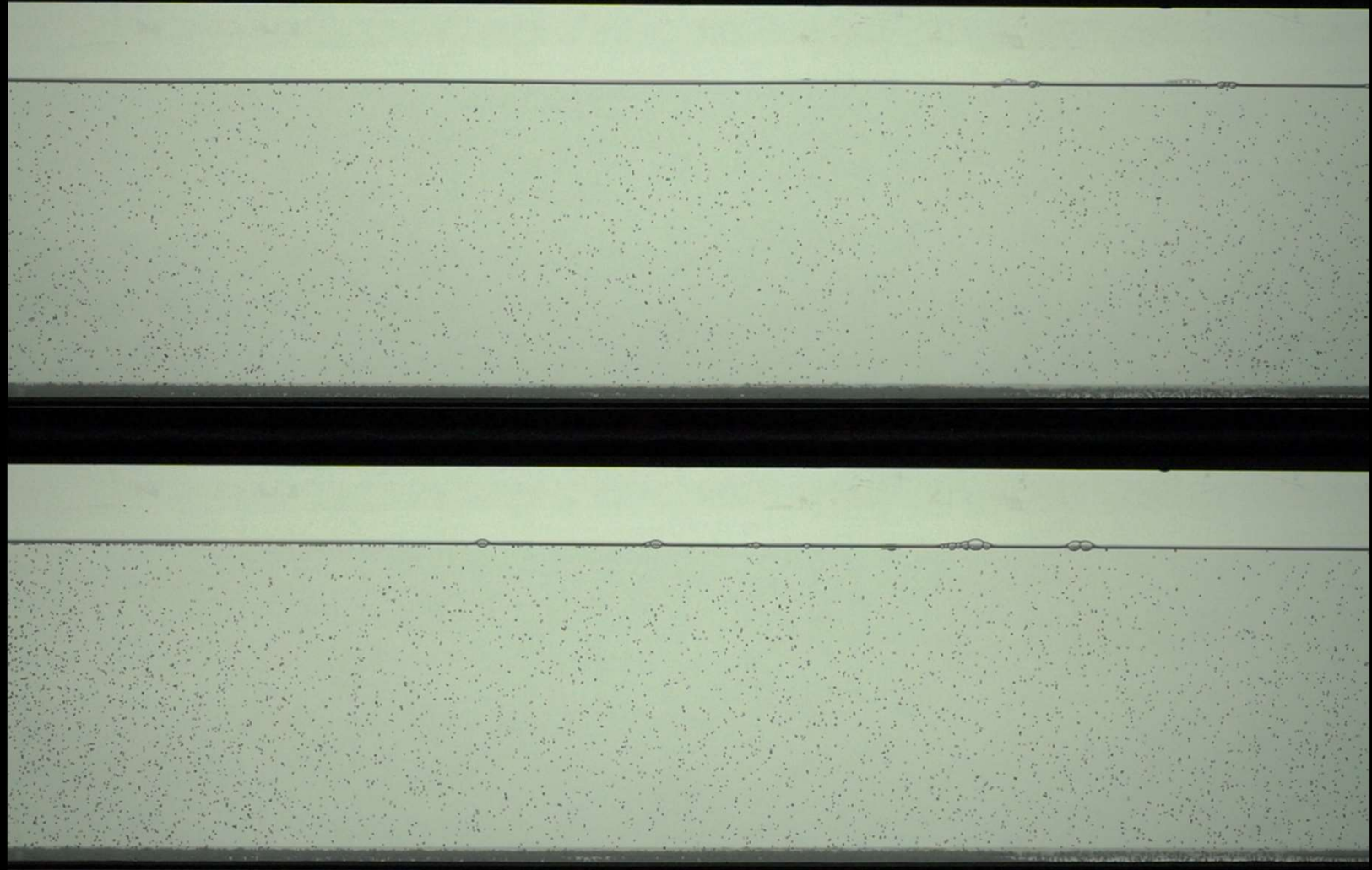
$$\Rightarrow \frac{\partial \hat{P}'}{\partial z} \Big|_{z=0} = \rho_0 g K \eta$$

$$\Rightarrow \frac{\partial^2 \eta}{\partial t^2} = -g K \eta$$

$$\Rightarrow \boxed{\omega^2 = g K} \text{ Dispersion relation}$$

# Surface gravity waves (SGWs)

$$\omega = \sqrt{gK}$$
$$P \sim \rho_0 g \eta e^{-Kz}$$



# SGWs

$$\omega = \pm\sqrt{gK}$$

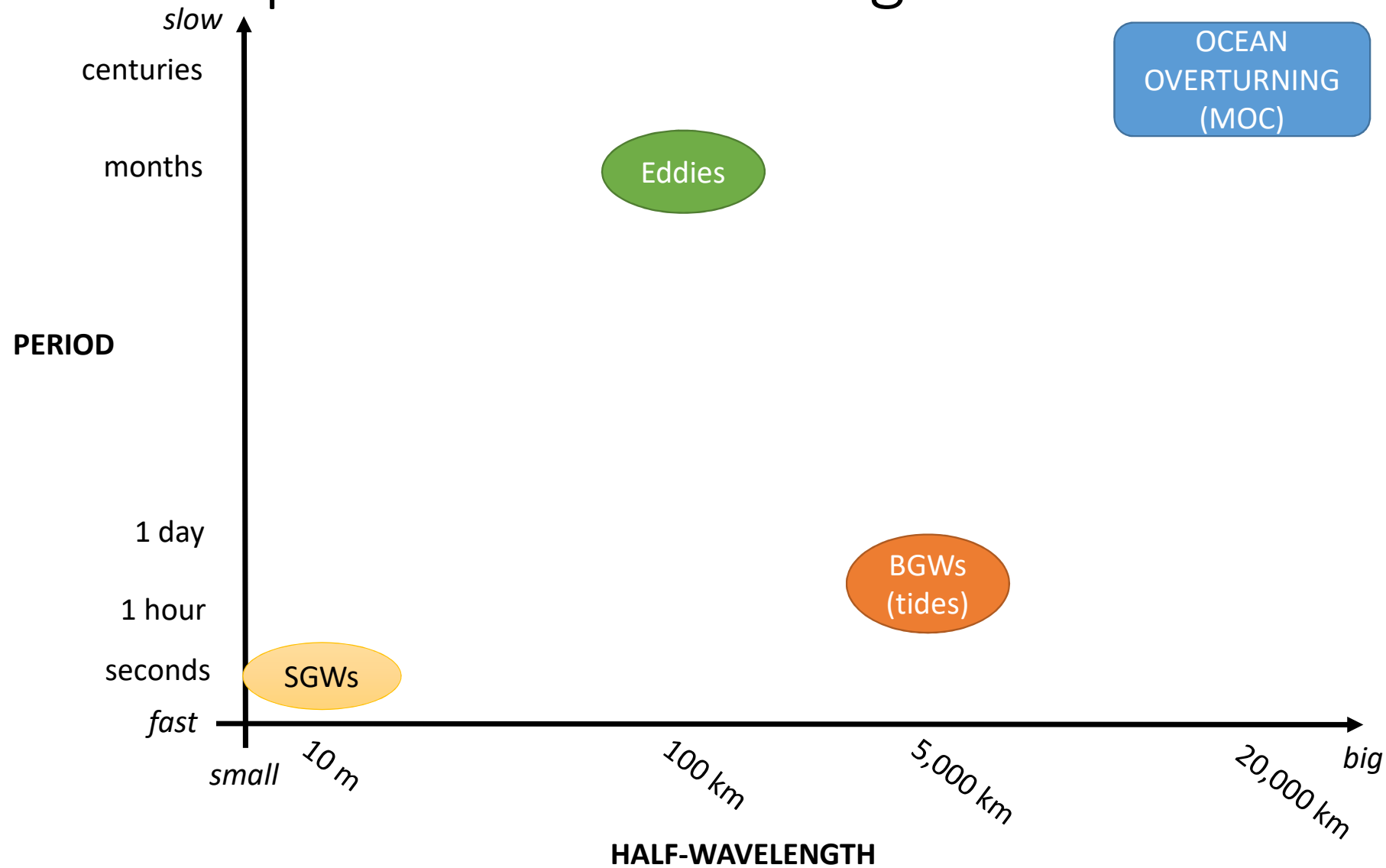
- Generated by perturbations at the surface (wind!)
- $\frac{2\pi}{\omega} \sim 3 \text{ seconds} \Rightarrow \frac{2\pi}{k} \sim 14 \text{ m}$
- Small spatial scales
- Non-hydrostatic = large vertical acceleration
- Shallow, surface intensified motion

# BGWs

$$\omega = \pm\sqrt{f^2 + K^2 gH}$$

- Generated by full depth/body force perturbations (tidal forces)
- $\frac{2\pi}{\omega} \sim 12.4 \text{ hrs} \Rightarrow \frac{2\pi}{k} \sim 10,000 \text{ km}$
- Large spatial scale
- Hydrostatic = slow vertical acceleration
- Full-depth barotropic motion

# Ocean space-time scale diagram



# Some references

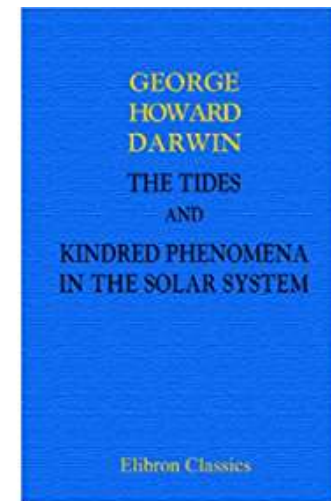
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# Spherical coordinates: water world

$$\nabla_h \cdot \mathbf{A} = \frac{1}{r \cos \theta} \frac{\partial}{\partial \phi} A_\phi + \frac{1}{r \cos \theta} \frac{\partial}{\partial \theta} A_\theta \cos \theta$$

$$\nabla_h A = \left( \frac{1}{r \cos \theta} \frac{\partial A}{\partial \phi}, \frac{1}{r} \frac{\partial A}{\partial \theta} \right)$$

$$f = 2\Omega \sin \theta$$

